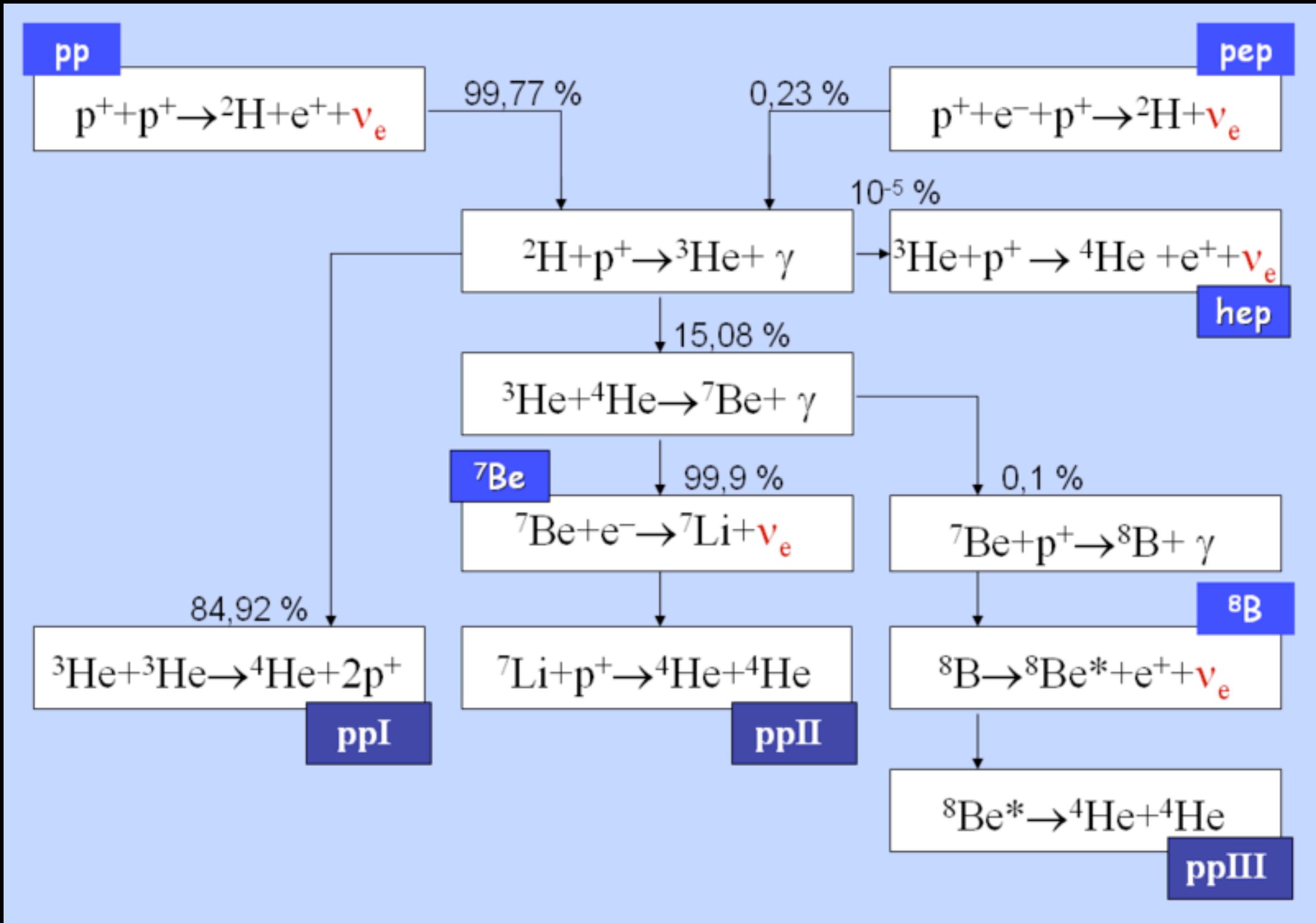


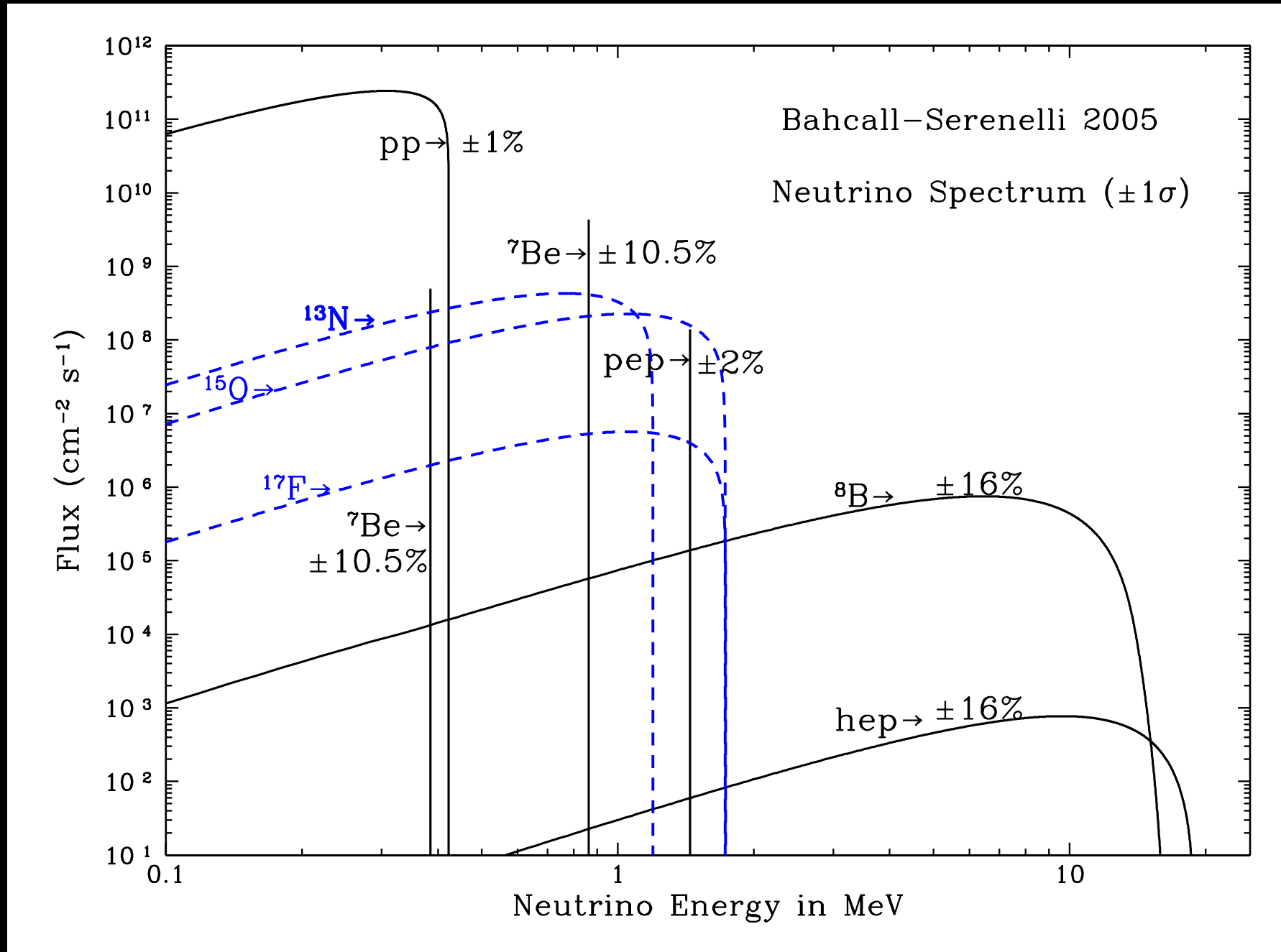
# The Solar Neutrino Problem

effective solar fusion reaction





$$\text{SNU} = 10^{-36} \text{ events/atom/s}$$



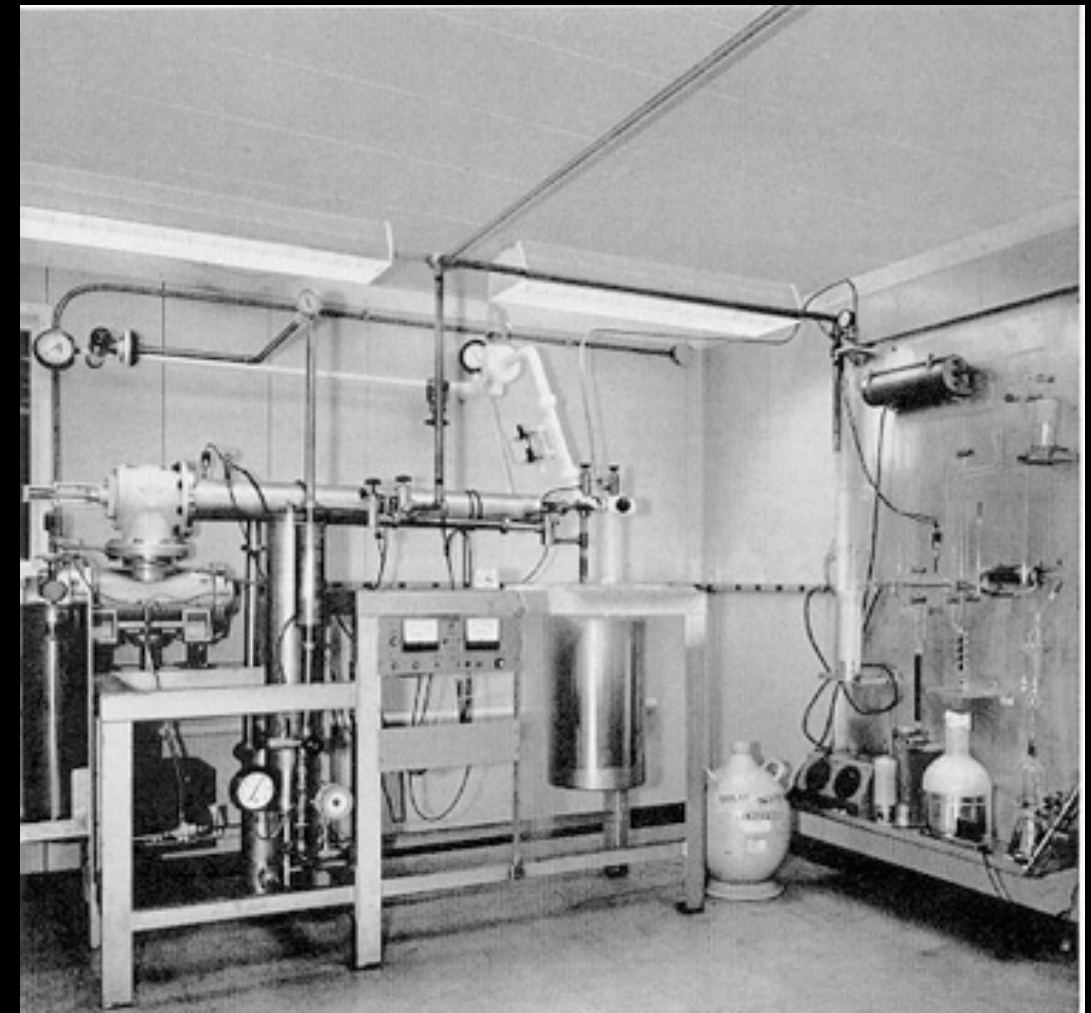
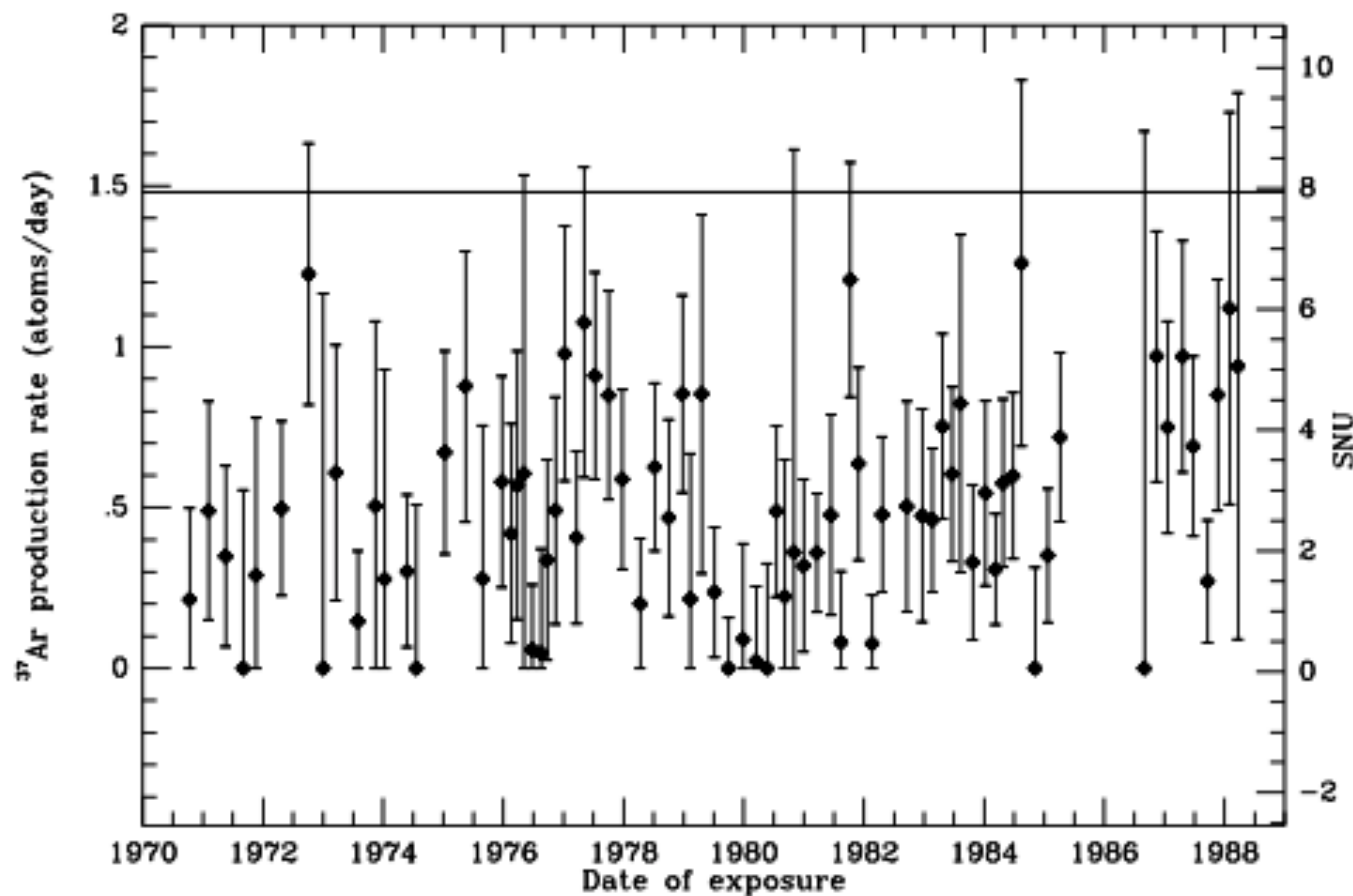
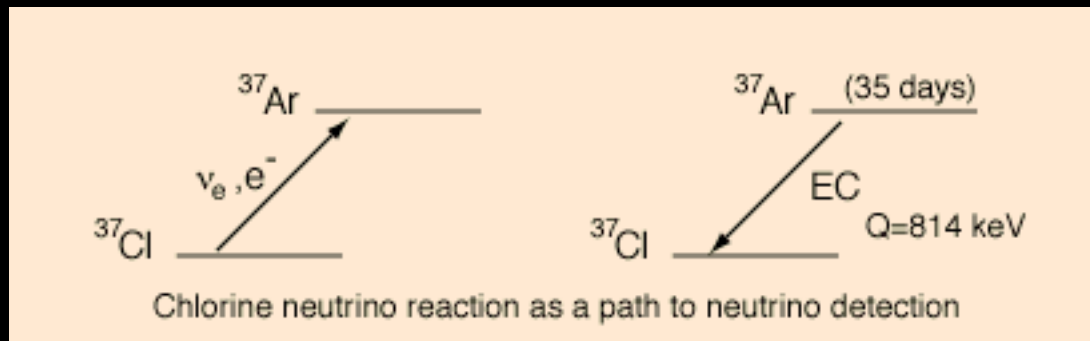
The predicted solar neutrino energy spectrum. The figure shows the energy spectrum of solar neutrinos predicted by the BP04 solar model. For continuum sources, the neutrino fluxes are given in number of neutrinos  $\text{cm}^{-2}\text{s}^{-1} \text{MeV}^{-1}$  at the Earth's surface. For line sources, the units are number of neutrinos  $\text{cm}^{-2}\text{s}^{-1}$ . Total theoretical uncertainties are shown for each source. To avoid complication in the figure, we have omitted the difficult-to-detect CNO neutrino fluxes.



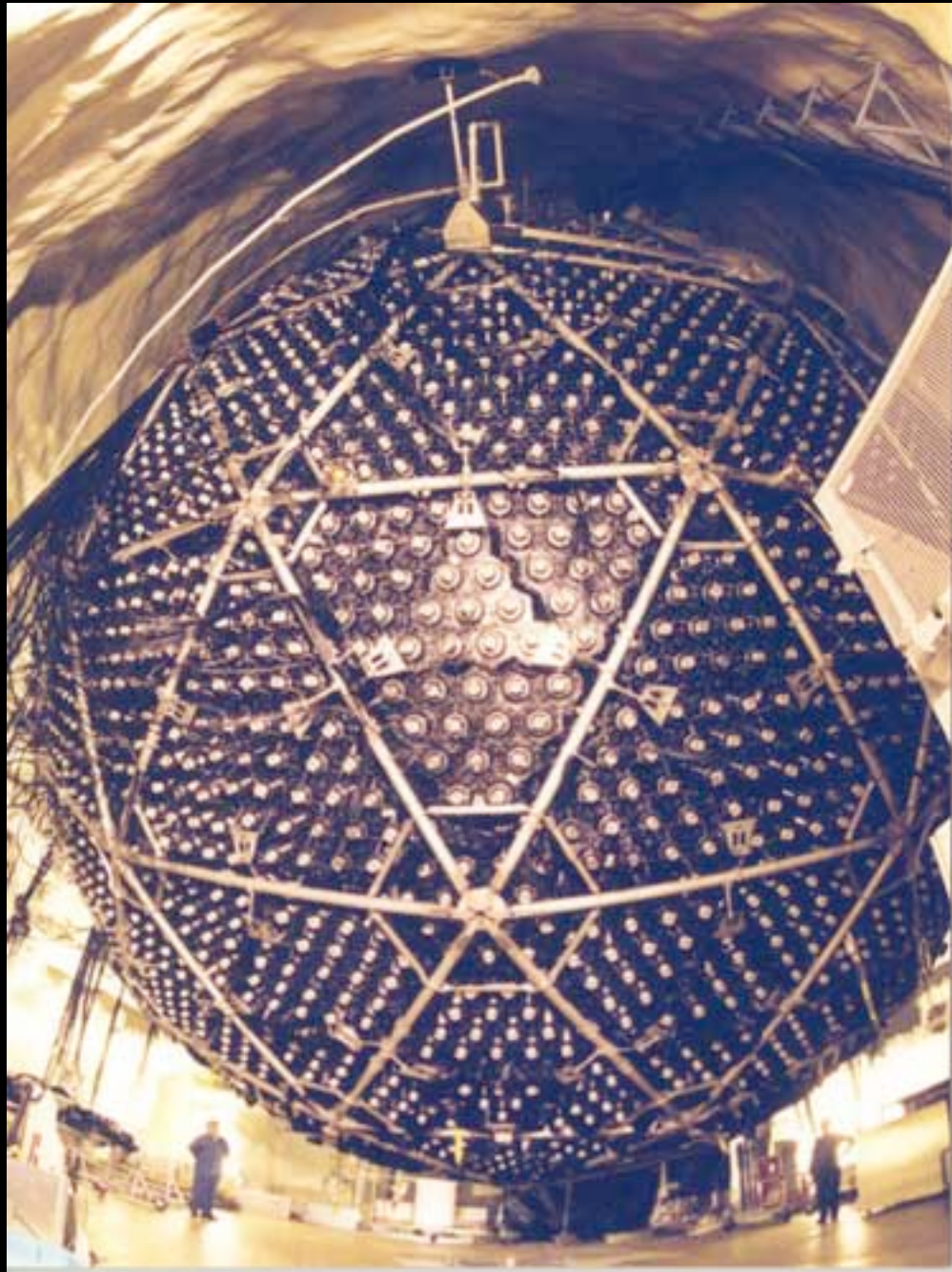
physics 330 Spring 2020

# 1964

Raymond Davis of Brookhaven National Laboratory constructed a neutrino detector 1.6 km underground in the Homestake Gold Mine in Lead, South Dakota. The detector consists of a 378,000 liter tank of perchloroethylene, which is further isolated by being submerged in water. Theoretical expectations were about one neutrino-chlorine interaction per day, but the measured solar neutrino events were about a third of that, raising serious questions about the abundance of solar neutrinos







$$\nu_e \rightarrow W^+ + e^-$$

$$\nu_\mu \rightarrow W^+ + \mu^-$$

$$\nu_\tau \rightarrow W^+ + \tau^-$$

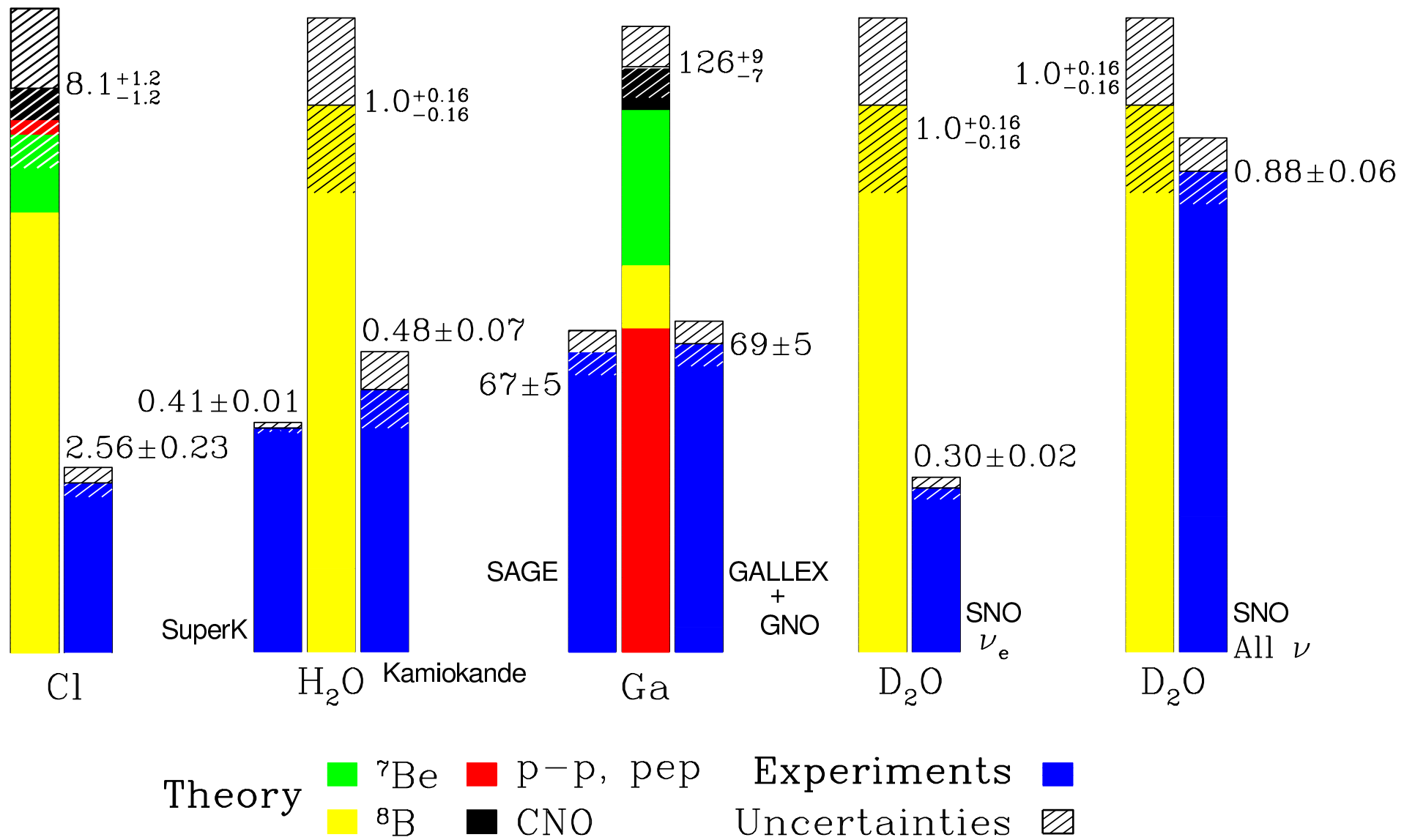
$$\nu \rightarrow Z + \nu$$

Sudbury Neutrino Observatory, a 12-meter sphere filled with heavy water surrounded by light detectors located two thousand meters below the ground in Sudbury, Ontario, Canada.

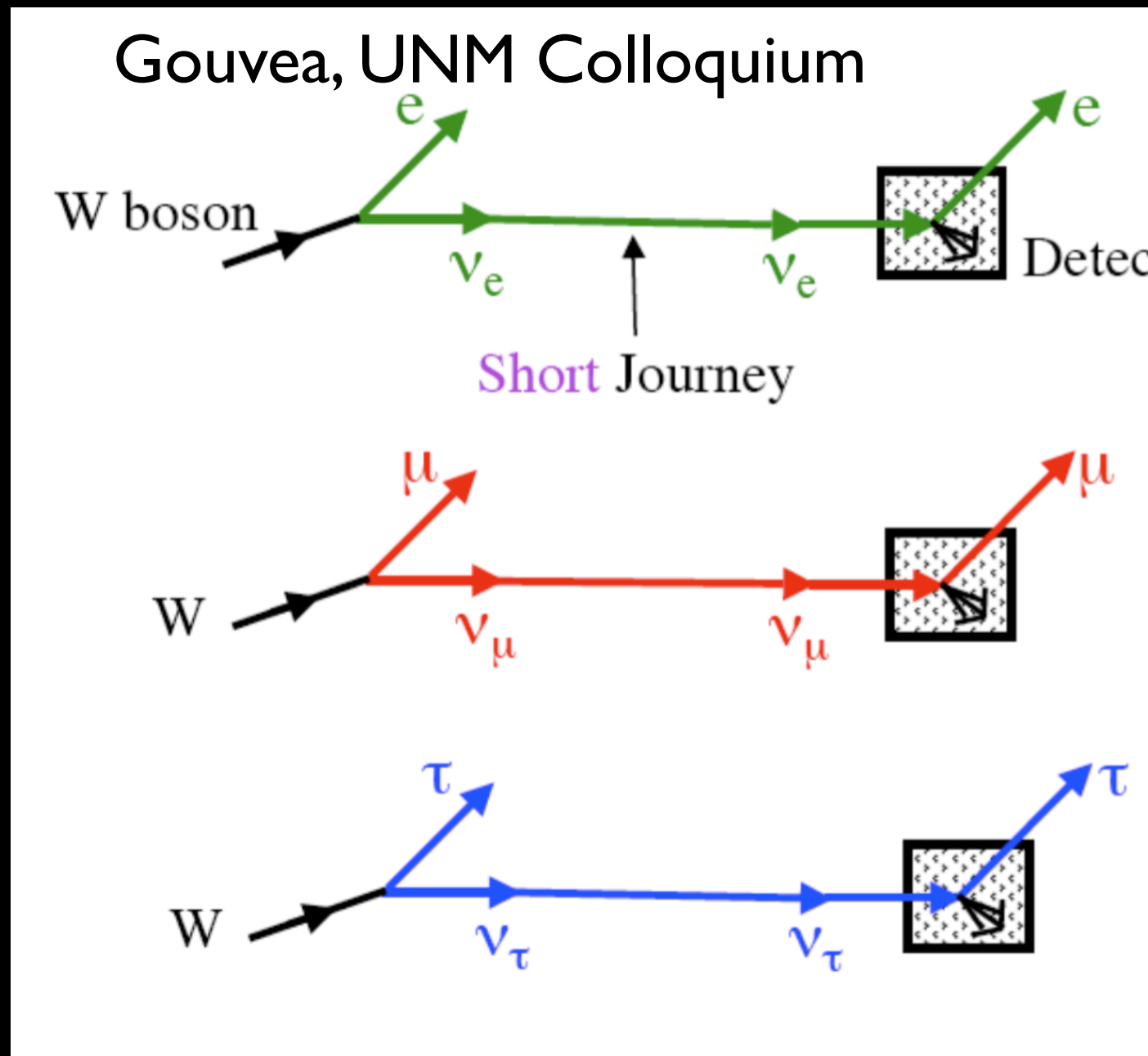
$$d + \nu \rightarrow p + n + \nu$$

# Total Rates: Standard Model vs. Experiment

Bahcall–Serenelli 2005 [BS05(OP)]



# Neutrino Mixing implies neutrino mass





## Neutrino Mixing

For simplicity, 2 massive neutrino state masses  $m_1, m_2$ .

$$\begin{pmatrix} \chi_e \\ \chi_\mu \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \chi_1 \\ \chi_2 \end{pmatrix}$$

orthogonal  $\chi_i^\dagger \chi_j = \delta_{ij}$

at  $t=0$ , produce  $\nu_e$  flavor

$$\Psi_e(0) = \cos\theta \chi_1 + \sin\theta \chi_2$$

Propagate as  
 $-iE_i + i\hbar$   
 $e$

ultra-relativistic

$$E_i = \sqrt{p^2 + m_i^2} \approx p + \frac{m_i^2}{2p}$$

Then there is a phase difference

$$\Psi_e(t) = \underbrace{e^{-iPt/\hbar}}_{\text{overall phase can be ignored}} \underbrace{e^{-im_1^2 t/\hbar}}_{-i\phi_{21} t} \left[ \cos\theta \chi_1 + e^{i\phi_{21} t} \sin\theta \chi_2 \right]$$

$$\phi_{21} = \frac{m_2^2 - m_1^2}{2p\hbar}$$

Probability to detect  $\nu_\mu$  flavor

$$P(e \rightarrow \mu) = \left| \chi_\mu^\dagger \Psi_e(t) \right|^2$$

neutrino mixing:

then

$$\chi_{\mu}^{+} = -\sin\theta \chi_{1}^{+} + \cos\theta \chi_{2}^{+}$$

using orthogonality

$$\chi_{\mu}^{+} \psi_{e}(t) = -\sin\theta \cos\theta + \sin\theta \cos\theta e^{-i\phi_{21}t}$$

$$= \underbrace{-\sin\theta \cos\theta}_{\frac{1}{2}\sin\left(\frac{\theta}{2}\right)} e^{-i\phi_{21}t/2} \underbrace{\begin{bmatrix} e^{i\phi_{21}t/2} & -i\phi_{21}t/2 \\ e^{-i\phi_{21}t/2} & \end{bmatrix}}_{2i \sin\left(\frac{\phi_{21}t}{2}\right)}$$

$$P(e \rightarrow \mu) = |\chi_{\mu}^{+} \psi_{e}(t)|^2$$
$$= \sin^2\left(\frac{\theta}{2}\right) \sin^2\left(\frac{\phi_{21}t}{2}\right)$$

with  $L = ct$ ,  $\phi \approx E$

$$P(e \rightarrow \mu) = \sin^2\left(\frac{\theta}{2}\right) \sin^2\left(\frac{\Delta m^2 L}{4E\hbar c}\right)$$

depends on square mass difference

$$\Delta m^2 = m_2^2 - m_1^2$$