

Bose-Einstein Condensate

March 23 2017

The Nobel Prize in Physics 2001 was awarded jointly to Eric A. Cornell, Wolfgang Ketterle and Carl E. Wieman "for the achievement of Bose-Einstein condensation in dilute gases of alkali atoms, and for early fundamental studies of the properties of the condensates".

A system of particles with integral spin obeys Bose-Einstein statistics.

$$f_{BE} = \frac{1}{Ae^{E/k_B T} - 1} = \frac{1}{e^{(E-\mu)/k_B T} - 1}$$

$\mu \equiv$ chemical potential $\mu < 0$ generally a function of temperature (T). Physically, $-\mu \equiv$ energy required to remove one molecule.

$$\frac{dn}{d^3p} = (2s + 1) \frac{1}{(2\pi\hbar)^3} f_{BE}(E)$$

The total number density n is fixed,

$$n = (2s + 1) \int_0^\infty \frac{4\pi p^2 dp}{(2\pi\hbar)^3} f_{BE}(E)$$

The non-relativistic energy is $E = p^2/2m$ so

$$n = (2s + 1) \int_0^\infty \frac{4\pi\sqrt{2mE}}{(2\pi\hbar)^3} m f_{BE}(E) dE$$

let $x = E/(k_B T)$

then

$$n = (2s+1) \left(\frac{2mk_B T}{\hbar^2} \right)^{3/2} \int_0^\infty \frac{\sqrt{x} dx}{\exp(x - \mu/(k_B T)) - 1}$$

Some atoms reach a critical temperature where $\mu(T_c) = 0$ so that

$$n = (2s + 1) \left(\frac{2mk_B T}{\hbar^2} \right)^{3/2} \int_0^\infty \frac{\sqrt{x} dx}{e^x - 1}$$

The value the dimensionless integral is $I = \frac{\sqrt{\pi}}{2}\zeta(3/2) = 2.32$ using the Zeta function $\zeta(3/2) = 26.12$

We can solve for T_c ,

$$T_c = \frac{2\pi\hbar^2}{mk_B} \left(\frac{n}{(2s+1)\zeta(3/2)} \right)^{2/3}$$

Below T_c , Einstein predicted that a fraction of the atoms will collapse into the same quantum state known as a Bose-Einstein condensate.

Below T_c

$$n = n_{cond} + \frac{2s + 1}{4\pi^2} (2.32) \left(\frac{2mk_B T}{\hbar^2} \right)^{3/2}$$

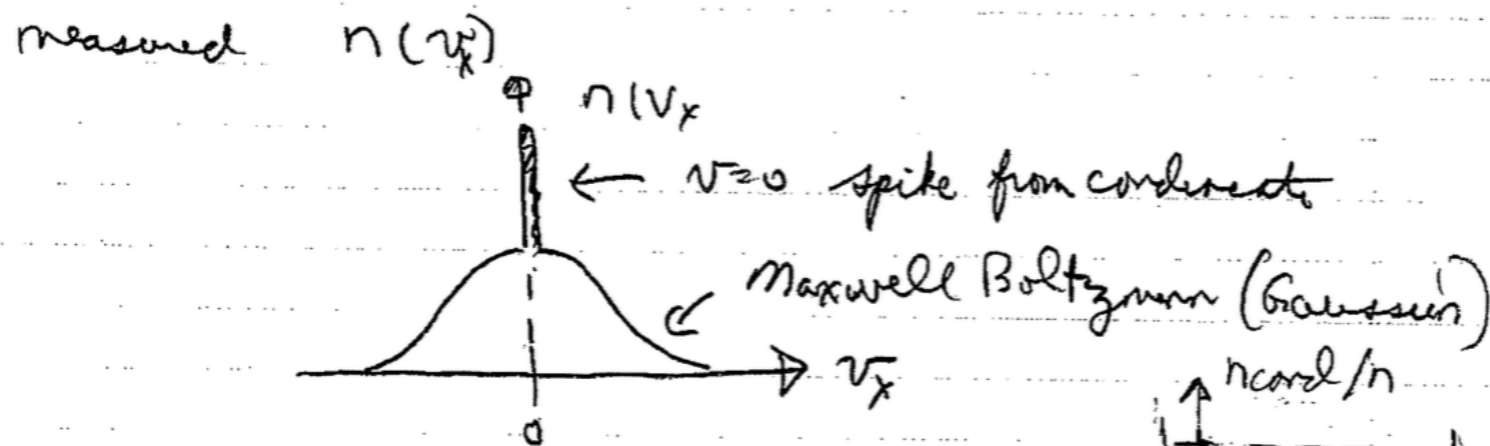
$$n = n_{cond} + n \left(\frac{T}{T_c} \right)^{3/2}$$

Giving for number density of the condensate
(n_{cond})

$$n_{cond} = n \left[1 - \left(\frac{T}{T_c} \right)^{3/2} \right]$$

1995, conclusively observed in dilute

gas of ^{87}Rb $T_c = 190 \times 10^{-9} \text{K}$



measured $\frac{n_{\text{cond}}(T)}{n} = 1 - \left(\frac{T}{T_c}\right)^{3/2}$

n_{cond}/n

T/T_c

Related phenomena are other types of

condensates (Bosons collapsed into ground state)

superfluidity, superconductivity.

In superfluid liquid helium, He atoms are interacting (chemical potential $\neq 0$), so that entire liquid is in ground state and exhibits zero viscosity.

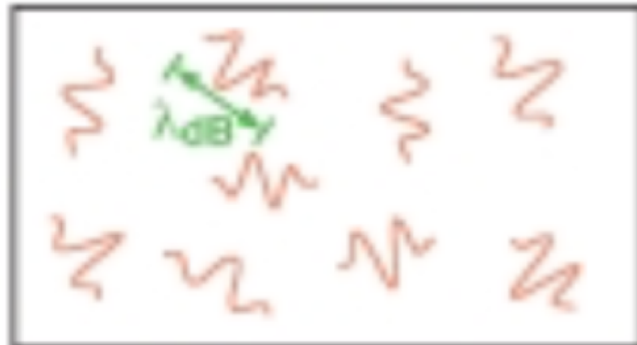
Atom is Boson if nucleons + electrons is even

$^{87}\text{Rb}(Z=37)$ is a Boson

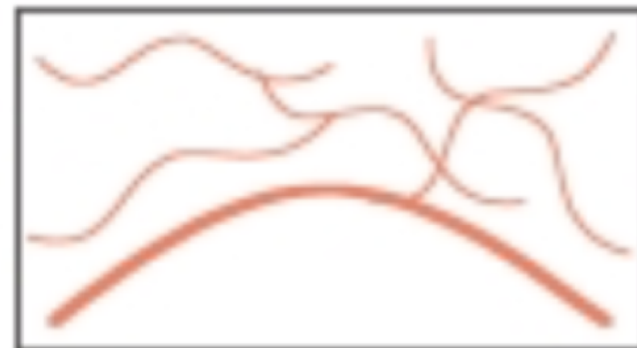
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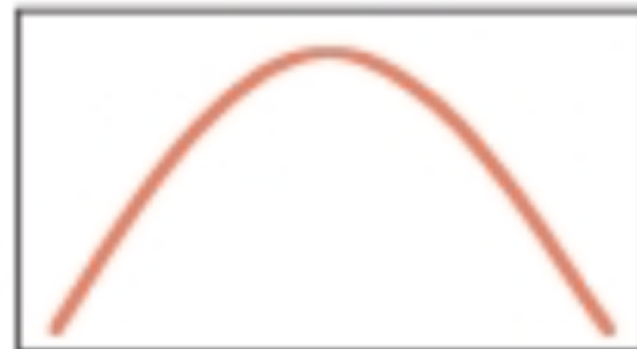
High Temperature T:
 thermal velocity v
 density d^{-3}
 "Billiard balls"



Low Temperature T:
 De Broglie wavelength
 $\lambda_{dB} = h/mv \propto T^{-1/2}$
 "Wave packets"



**T = T_{crit}:
 Bose-Einstein
 Condensation**
 $\lambda_{dB} \sim d$
 "Matter wave overlap"

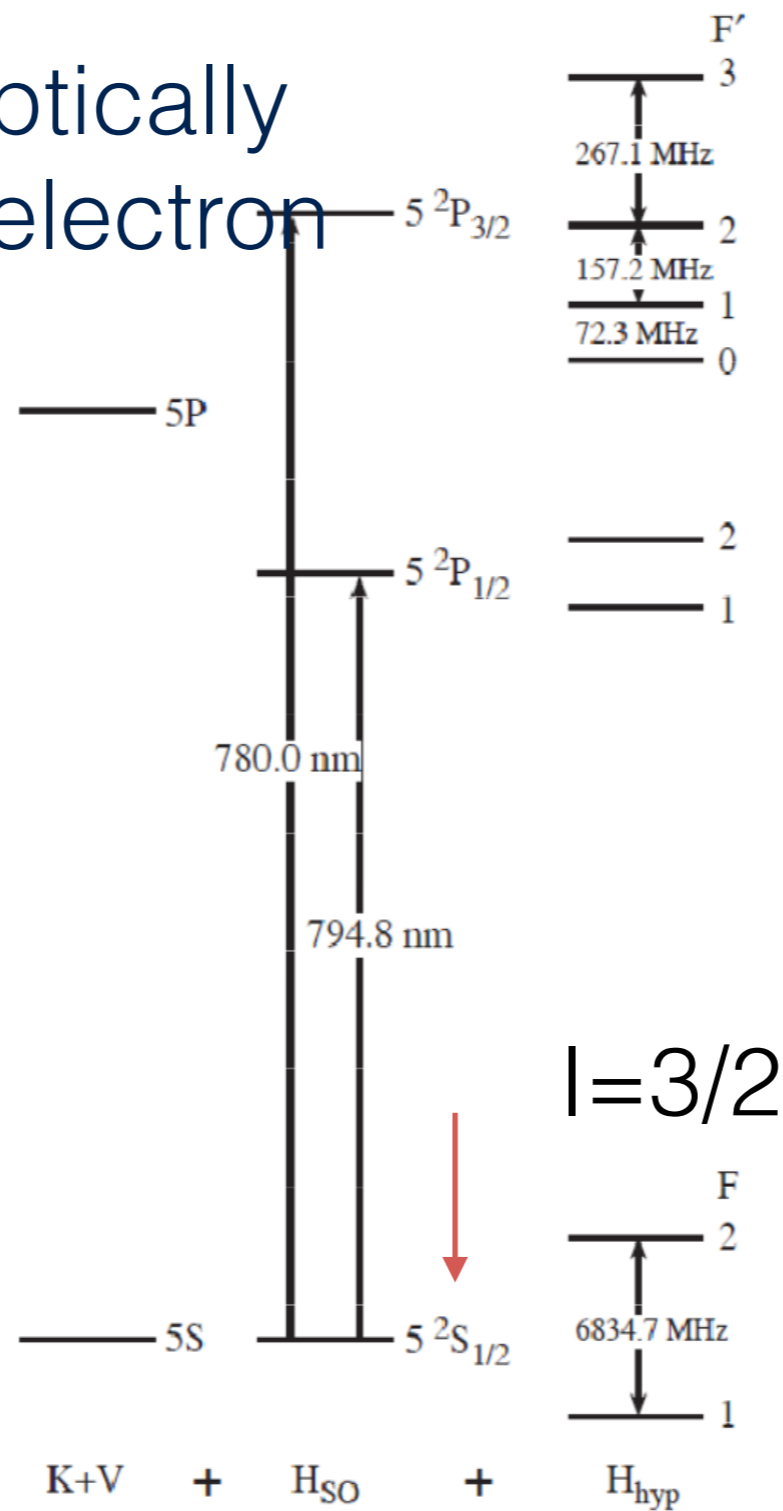


**T = 0:
 Pure Bose
 condensate**
 "Giant matter wave"

Figure 2. Criterion for Bose-Einstein condensation. At high temperatures, a weakly interacting gas can be treated as a system of "billiard balls." In a simplified quantum description, the atoms can be regarded as wavepackets with an extension of their de Broglie wavelength λ_{dB} . At the BEC transition temperature, λ_{dB} becomes comparable to the distance between atoms, and a Bose condensate forms. As the temperature approaches zero, the thermal cloud disappears, leaving a pure Bose condensate.

37 Rb rubidium : [Kr] 5s¹

one optically active electron



$m_F = 1, 0, -1$

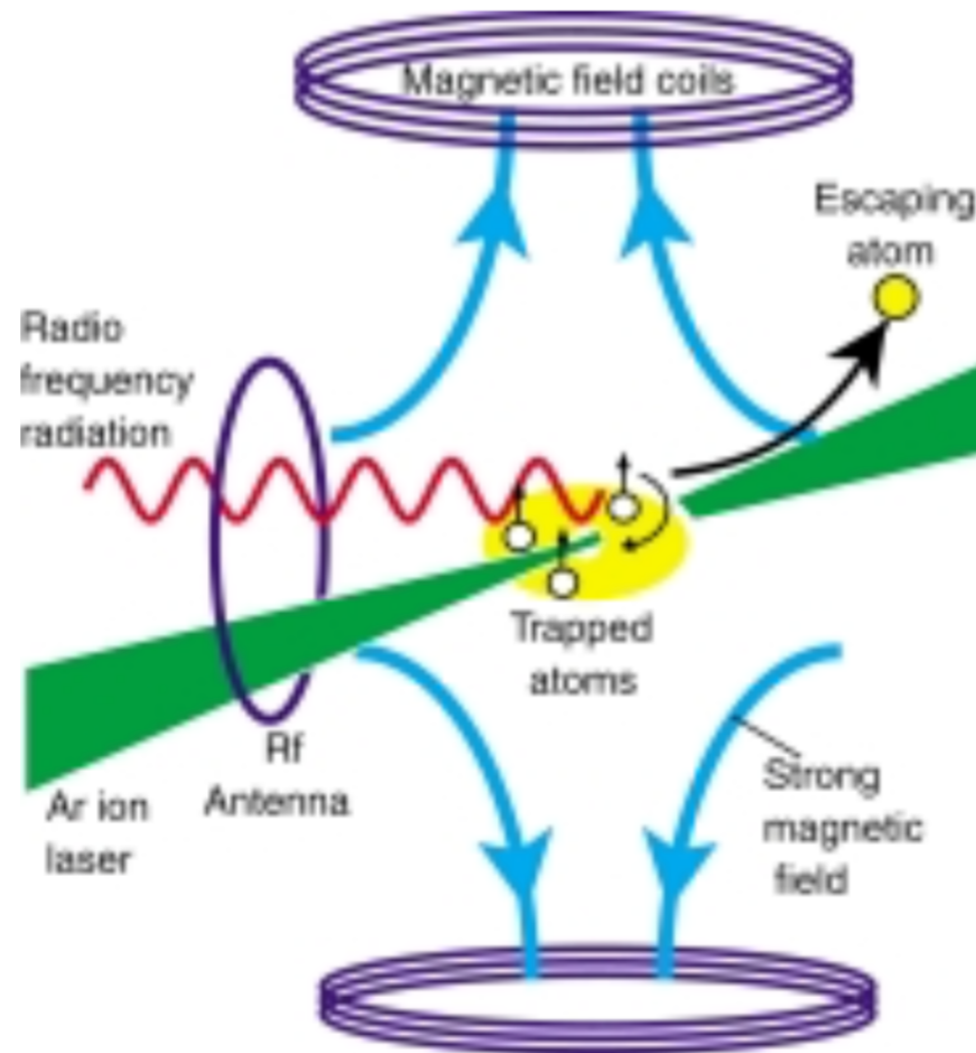


Figure 3. Experimental setup for cooling atoms to Bose-Einstein condensation. Sodium atoms are trapped by a strong magnetic field, generated by two coils. In the center, the magnetic field vanishes, which allows the atoms to spin-flip and escape. Therefore, the atoms are kept away from the center of the trap by a strong (3.5 W) argon ion laser beam (“optical plug”), which exerts a repulsive force on the atoms. Evaporative cooling is controlled by radio-frequency radiation from an antenna. The rf selectively flips the spins of the most energetic atoms. The remaining atoms rethermalize (at a lower temperature) by collisions among themselves. Evaporative cooling is forced by lowering the rf frequency.

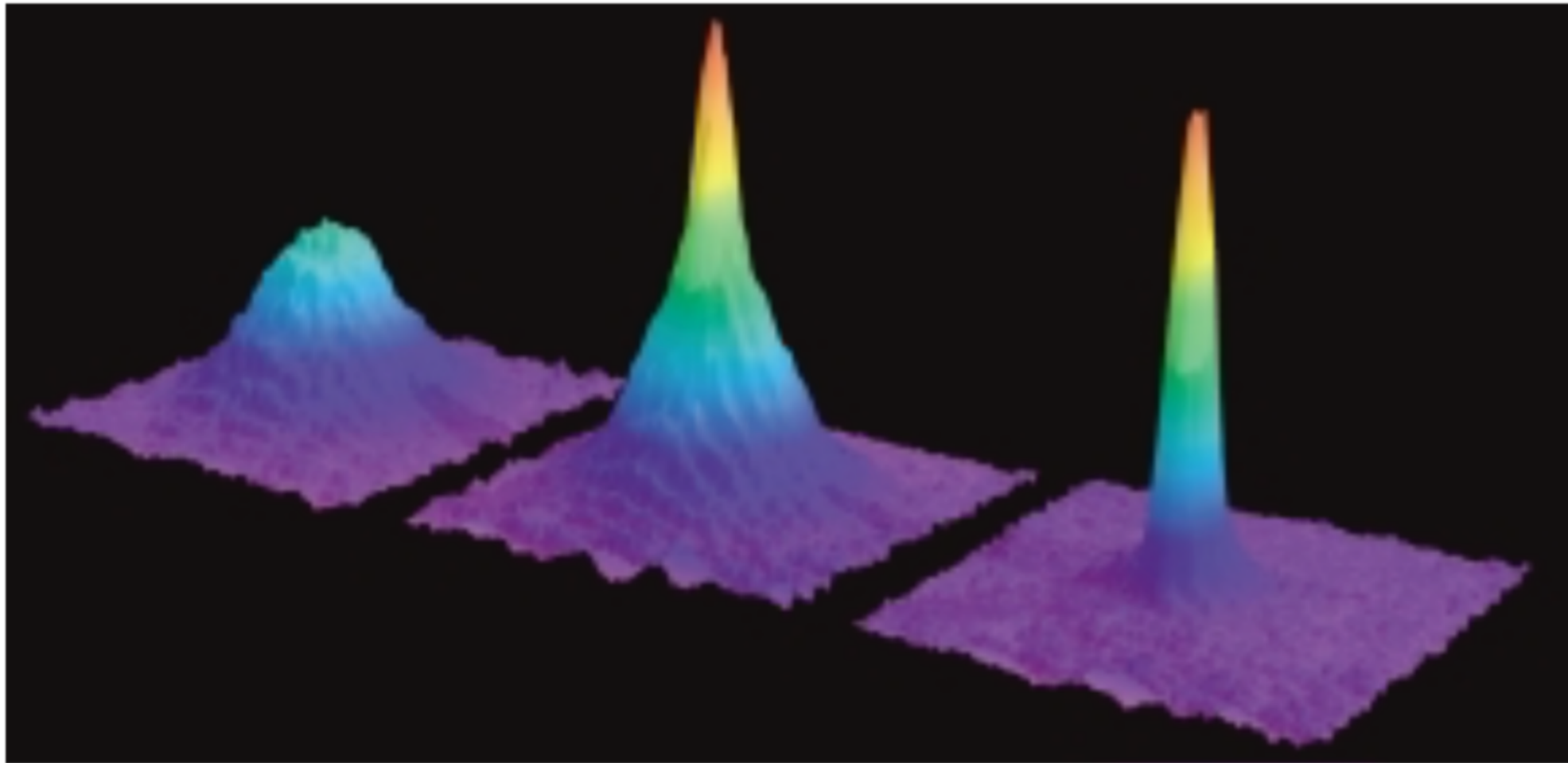
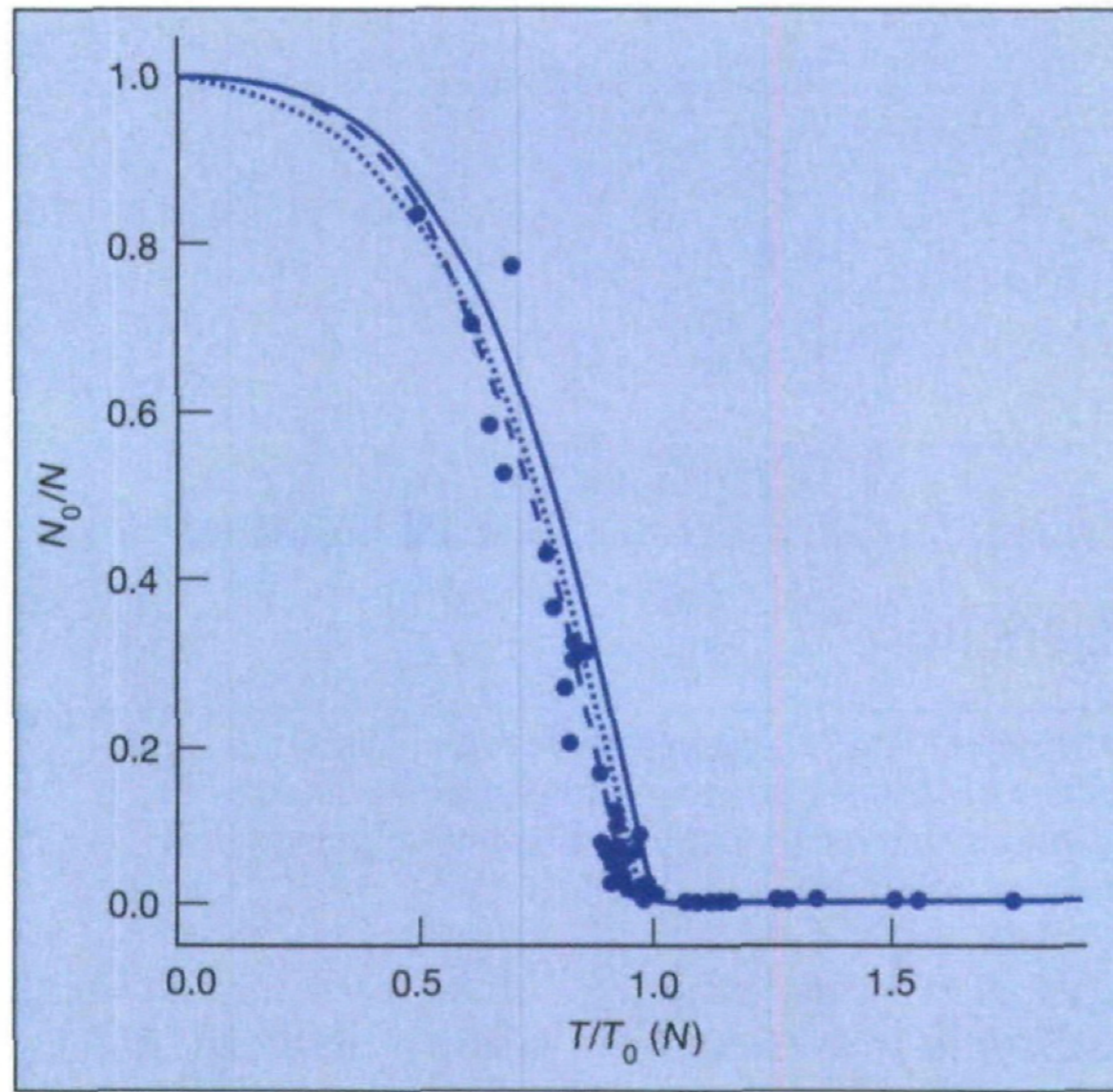


Figure 7. Observation of Bose-Einstein condensation by absorption imaging. Shown is absorption vs. two spatial dimensions. The Bose-Einstein condensate is characterized by its slow expansion observed after 6 ms time-of-flight. The left picture shows an expanding cloud cooled to just above the transition point; middle: just after the condensate appeared; right: after further evaporative cooling has left an almost pure condensate. The total number of atoms at the phase transition is about 7×10^5 , the temperature at the transition point is $2 \mu\text{K}$.



5 Condensate fraction, N_0/N , measured as a function of scaled temperature, T/T_0 , at Boulder. T_0 is the predicted critical temperature for an ideal gas in a harmonic trap in the thermodynamic limit (many atoms). The solid curve is the predicted dependence in the thermodynamic limit, and the dotted curve includes a correction for the finite number of atoms (4000) in the condensate. To the uncertainties in the data, the measurement is consistent with the theory for non-interacting bosons. The dashed line is a best fit to the experimental data.