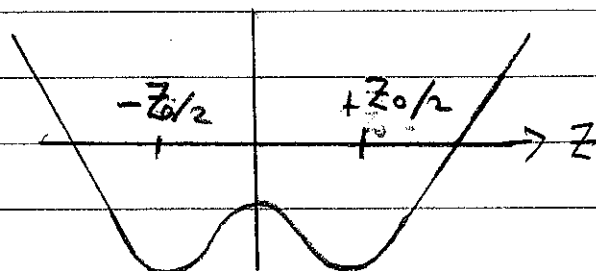


Lecture 16: Covalent Bond

Approximate proton separation as fixed.
Potential seen by electron is double-well:



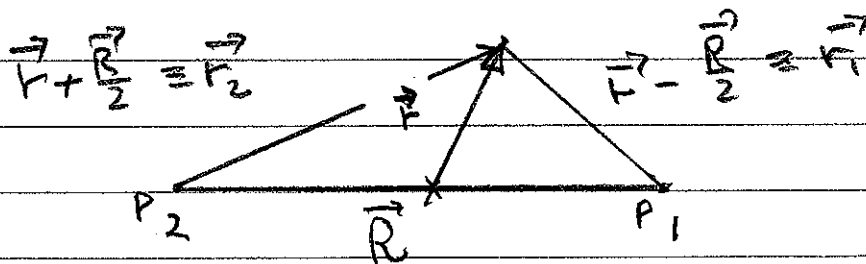
$$V(z) = \frac{m\omega^2}{2} \left(|z| - \frac{z_0}{2} \right)^2$$

Constant proton separation z_0 justified by
slow proton oscillations

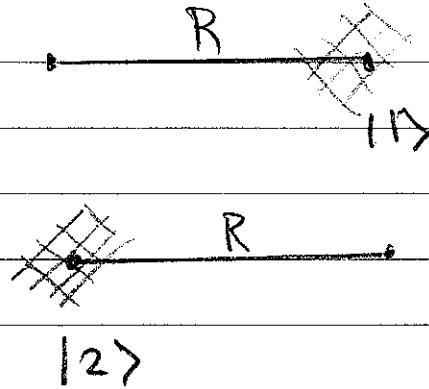
$$\frac{1}{2} m \omega^2 = \frac{1}{2} \left(\frac{m}{m_p} \right) m_p \omega^2$$

So expect $\omega_p \sim \sqrt{\frac{m}{m_p}} \omega$

Coordinate System:



perturbative approach: $H + P$



Exact
solution for
 $R \gg a_0$

$$\begin{aligned} \langle \vec{r}_1 | 1 \rangle &= \frac{1}{\sqrt{\pi a_0^3}} e^{-r_1/a_0} \\ \langle \vec{r}_2 | 2 \rangle &= \frac{1}{\sqrt{\pi a_0^3}} e^{-r_2/a_0} \end{aligned} \quad \left. \vphantom{\begin{aligned} \langle \vec{r}_1 | 1 \rangle &= \frac{1}{\sqrt{\pi a_0^3}} e^{-r_1/a_0} \\ \langle \vec{r}_2 | 2 \rangle &= \frac{1}{\sqrt{\pi a_0^3}} e^{-r_2/a_0} \end{aligned}} \right\} \begin{array}{l} \text{degenerate} \\ \text{unperturbed} \\ \text{states} \end{array}$$

$$\hat{H} = \frac{\vec{p}^2}{2m} - \frac{e^2}{r_1} - \frac{e^2}{r_2} + \frac{e^2}{R}$$

H matrix:

$$\langle 1 | \hat{H} | 1 \rangle = \langle 2 | \hat{H} | 2 \rangle = H_{11} = H_{22} = E_0$$

$$\langle 1 | \hat{H} | 2 \rangle = \langle 2 | \hat{H} | 1 \rangle = H_{12} = H_{21} = -A$$

Where $-A$ is energy due to tunneling
across potential barrier separating states
 $|1\rangle, |2\rangle$.

Diagonalizing $[H] = \begin{pmatrix} E_0 & -A \\ -A & E_0 \end{pmatrix}$

solutions:

$$E_I = E_0 - A, \quad |I\rangle = \frac{1}{\sqrt{2}}(|1\rangle + |2\rangle) \quad \text{binding}$$

$$E_{II} = E_0 + A, \quad |II\rangle = \frac{1}{\sqrt{2}}(|1\rangle - |2\rangle) \quad \text{non-binding}$$

we have assumed $\langle 1|2\rangle = 0$ which is not correct.
Introduce overlap

$\Delta = \langle 1|2\rangle = \langle 2|1\rangle$ real, dimensionless
 get (see HW)

$$E_I = \frac{1}{1+\Delta} (E_0 - A); \quad |I\rangle = \frac{1}{\sqrt{2+2\Delta}} (|1\rangle + |2\rangle)$$

$$E_{II} = \frac{1}{1-\Delta} (E_0 + A); \quad |II\rangle = \frac{1}{\sqrt{2-2\Delta}} (|1\rangle - |2\rangle)$$

Δ further lowers the ground state energy.
 Following Pauling and Wilson,

$$J \equiv \int d^3r_1 \frac{e^2}{r_2} |\langle n|1_1\rangle|^2 \propto \frac{e^2}{a_0}$$

with $E_1 \equiv \text{H-atom ground state energy}$

$$\langle 1|\hat{H}|1\rangle = \langle 2|\hat{H}|2\rangle = E_1 - J + \frac{e^2}{R} = E_0$$

and

$$K \equiv \int d^3 r_1 \langle 2 | r_1 \rangle \frac{e^2}{r_2} \langle r_1 | 1 \rangle$$

off diagonal (tunneling) terms are

$$\langle 2 | \hat{H} | 1 \rangle = \Delta \cdot E_1 - K + \Delta \frac{e^2}{R} = -A$$

$$\text{giving } A = K - \Delta \left(E_1 + \frac{e^2}{R} \right)$$

then in terms of J, K

$$\begin{aligned} E_I &= \frac{1}{1+\Delta} \left[E_1 + J + \frac{e^2}{R} - K + \Delta \left(E_1 + \frac{e^2}{R} \right) \right] \\ &= E_1 + \frac{e^2}{R} - \left(\frac{J+K}{1+\Delta} \right) \end{aligned}$$

$$E_{II} = E_1 + \frac{e^2}{R} - \left(\frac{J-K}{1-\Delta} \right)$$

for completeness, integrals are ($D \equiv \frac{R}{a_0}$)

$$\Delta = e^{-D} \left(1 + D + \frac{1}{3} D^2 \right)$$

$$J = -\frac{e^2}{a_0} \left[-\frac{1}{D} + e^{-2D} \left(1 + \frac{1}{D} \right) \right]$$

$$K = \frac{e^2}{a_0} e^{-D} \left(1 + \Delta \right)$$

$$\text{where } \frac{e^2}{a_0} = \hbar c \alpha \left(\frac{mc \alpha}{\hbar} \right) = mc^2 \alpha^2$$

E_I is then a function of R and has a minimum @ $R = 0.13 \text{ nm}$ and value -15.4 eV

binding or dissociation = pulling the molecule apart into proton and hydrogen

binding energy is $-E_{\text{bind}} = E_I - E_1$

after dissociation, electron is bound to one of the protons

$$E_{\text{bind}} = (-13.6 \text{ eV}) - (-15.4 \text{ eV}) = 1.8 \text{ eV}$$

Exp value $R = 0.106 \text{ nm}$ $E_{\text{bind}} = 2.8 \text{ eV}$

We have assumed wave functions are undistorted by presence of other proton.

Covalent bond is due to tunneling between degenerate states.