

Lec 9: Simple Harmonic Oscillator

Importance: Near potential minimum, almost all potentials are harmonic.

$$V(x) = V(x_0) + (x-x_0)V'(x_0) + \frac{1}{2}(x-x_0)^2 V''(x_0) + \dots$$

0'' for $x_0 = \text{minimum}$

as long as $V''(x_0) \neq 0$

$$V(x) - V(x_0) = \frac{1}{2}(x-x_0)^2 V''(x_0)$$

$$V''(x_0) = k = m\omega^2 \quad \omega = \text{characteristic frequency} \times 2\pi$$

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2$$

define dimensionless variables

$$\hat{q} = \frac{\hat{p}}{\sqrt{m\hbar\omega}} \quad \hat{y} = \hat{x} \sqrt{\frac{m\omega}{\hbar}}$$

turns out, classical turning point of ground state is $x_c = \sqrt{\frac{\hbar}{m\omega}}$ ($E_0 = \frac{\hbar\omega}{2}$)

$$\hat{y} = \frac{\hat{x}}{x_c}; \quad \hat{q} = \frac{x_c}{\hbar} \hat{p}$$

dimensionless energy $\varepsilon = \frac{E}{E_0} = \frac{2E}{\hbar\omega}$

note - different from Shankar

then $\hat{H} = \frac{1}{2} \hbar \omega (\hat{\xi}^2 + \hat{y}^2)$

$$[\hat{x}, \hat{p}] = i\hbar \Rightarrow [\hat{y}, \hat{\xi}] = i$$

implying in $|y\rangle$ basis $\hat{\xi} = \frac{1}{i} \frac{\partial}{\partial y}$

time independent Schrodinger ψ

$$\psi'' + (\epsilon - y^2)\psi = 0 \quad (\psi' \equiv \frac{\partial}{\partial y} \psi)$$

standard solution $|y| \rightarrow \infty$

$$\psi'' = y^2 \psi \quad \psi \rightarrow e^{-y^2/2}$$

let $\psi = u(y) e^{-y^2/2}$ $|y| \rightarrow \infty$

$$u'' - 2u u' + (\epsilon - 1)u = 0$$

power series solution $u(y) = \sum_{k=0}^{\infty} C_k y^k$

get recursion relation.

require $\psi(y) \rightarrow 0$ $|y| \rightarrow \infty$

leads to termination of recursion relation and quantization of E .

$$E = \hbar \omega \left(\frac{1}{2} + n \right) \quad n = 0, 1, 2, 3$$

The polynomials are standard Hermite polynomials

$$\Psi_n(x) = \left(\frac{m\omega}{\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n\left(\sqrt{\frac{m\omega}{\hbar}} x\right) e^{-\frac{m\omega x^2}{2\hbar}}$$

$$H_0(y) = 1$$

$$H_1(y) = 2y \quad (\text{look them up})$$

$$H_2(y) = 4y^2 - 2$$

Ground state is Gaussian

generating function: $H_n(y) = (-1)^n e^{y^2} \frac{\partial^n}{\partial y^n} e^{-y^2}$

series expansion: $e^{-y^2 + 2zy} = \sum_{n=0}^{\infty} \frac{z^n}{n!} H_n(y)$

Algebraic Solution (Dirac) "factorize"

$$\hat{a} \equiv \frac{1}{\sqrt{2}} (\hat{y} + i\hat{p})$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}} (\hat{y} - i\hat{p})$$

$$[\hat{y}, \hat{p}] = i \Rightarrow [\hat{a}, \hat{a}^\dagger] = 1$$

$$\text{then } \hat{a}^\dagger \hat{a} = \frac{1}{2} (\hat{y} - i\hat{p})(\hat{y} + i\hat{p}) = \frac{1}{2} \{ \hat{y}^2 + \hat{p}^2 + i[\hat{y}, \hat{p}] \}$$

$$= \frac{1}{2} (\hat{y}^2 + \hat{p}^2) - \frac{1}{2}$$

$$\hat{H} = \hbar\omega \left(\hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$$

$a^\dagger a$ is called the number operator \hat{n} .
define \hat{n} basis

$$\hat{n}|n\rangle = n|n\rangle$$

$$\hat{H}|n\rangle = \hbar\omega\left(\hat{n} + \frac{1}{2}\right)|n\rangle = \hbar\omega\left(n + \frac{1}{2}\right)|n\rangle$$

So far, all we know is n is real

$$\hat{n}^\dagger = (a^\dagger a)^\dagger = a^\dagger a$$

Consider

$$\begin{aligned} [\hat{n}, \hat{a}] &= a^\dagger a a - a a^\dagger a \\ &= a^\dagger a a - \overbrace{a a^\dagger} + a \\ &= a^\dagger a a - a^\dagger a a - a = -a \end{aligned}$$

also

$$[\hat{n}, \hat{a}^\dagger]^\dagger = [\hat{a}^\dagger, \hat{n}] = -a^\dagger$$

$$[\hat{n}, \hat{a}^\dagger] = +a^\dagger$$

the state $a^\dagger|n\rangle$... Eigenvalue is

$$\hat{n}(a^\dagger|n\rangle) = (a^\dagger n + a^\dagger)|n\rangle = (n+1)(a^\dagger|n\rangle)$$

$$\text{So } a^\dagger|n\rangle = C_+|n+1\rangle$$

$$\text{Similarly } a|n\rangle = C_-|n-1\rangle$$

C_\pm normalization constants.

taking hermitian conjugate

$$\langle n | a = c_+^* \langle n |$$

then

$$\langle n | a a^\dagger | n \rangle = |c_+|^2 \langle n | n \rangle$$

properly normalize all eigenstates $\langle n | n \rangle = 1$

$$\langle n | a a^\dagger | n \rangle = |c_+|^2$$

$$\langle n | (a^\dagger a + 1) | n \rangle = |c_+|^2$$

$$n + 1 = |c_+|^2 \quad \text{can choose phase} = 1 \text{ so}$$

$$c_+ = \sqrt{n+1}$$

similarly, $c_- = \sqrt{n}$.

then $|n+1\rangle = \frac{a^\dagger}{\sqrt{n+1}} |n\rangle$ properly normalized

$$|n\rangle = \frac{(a^\dagger)^n}{\sqrt{n!}} |0\rangle \quad |0\rangle \text{ ground state}$$

must exist ground state so that

$$a |0\rangle = 0$$

from $a |n\rangle = \sqrt{n} |n-1\rangle$ we see
eigenvalue of ground state is number 0.
That makes all eigenvalues integers.

$$n = 0, 1, 2, \dots$$

In y representation:

$$\langle y | n \rangle = \frac{1}{\sqrt{2^n n!}} \left(y - \frac{d}{dy} \right)^n \langle y | 0 \rangle$$

From $a | 0 \rangle = 0$ we get

$$0 = \langle y | \hat{a} | 0 \rangle = \frac{1}{\sqrt{2}} \left(y + \frac{d}{dy} \right) \langle y | 0 \rangle$$

$$\frac{d}{dy} \langle y | 0 \rangle = -y \langle y | 0 \rangle$$

simple, first order equation

$$\langle y | 0 \rangle = \frac{1}{\pi^{1/4}} e^{-y^2/2}$$

Generating formula for Hermite polynomials:

$$\left(y - \frac{d}{dy} \right)^n e^{-y^2/2} = H_n(y) e^{-y^2/2}$$

Cosmological Constant Problem

Each frequency mode of EM field is harmonic oscillator, with photons

$$a_{\vec{k}, \lambda} |0\rangle = |\vec{k}, \lambda\rangle \quad \lambda \text{ polarization}$$

Zero-point harmonic oscillator energy implies that each frequency mode has vacuum energy $\frac{1}{2} \hbar \omega = \frac{1}{2} \hbar c k$. Summing up gives

Vacuum energy density

$$\langle \rho_{EM} \rangle = \int_0^\Lambda \frac{4\pi k^2 dk}{(2\pi)^3} \left(\frac{\hbar}{2} \right) = \frac{\Lambda^4}{16\pi^2}$$

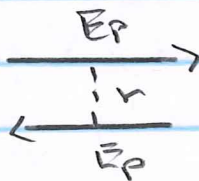
Where formally divergent integral is cut off at Planck scale where quantum field theory must break down,

$$\Lambda = \frac{1}{\sqrt{8\pi G}} \approx 10^{19} \text{ GeV}$$

One way to think about this scale: proton Coulomb repulsion compared to gravitational attraction.

Gravity couples to energy ($E_p = m_p c^2 \gamma$)

2 protons with ultra-high energy



proton attracted as $\frac{E_p}{c^2}$. Equate Coulomb repulsion to gravitational attraction:

$$\alpha \frac{\hbar c}{r} = \frac{G}{r} \left(\frac{E_p}{c^2} \right)^2$$

ignore the factor α , $\hbar c = G \left(\frac{E_p}{c^2} \right)^2$

$$E_p = c^2 \sqrt{\frac{\hbar c}{G}} \quad \text{so} \quad \Lambda \propto \frac{1}{\sqrt{G}}$$

in suitable units $\hbar=1, c=1$ $\Lambda = \frac{1}{\sqrt{8\pi G}}$

$$\text{then } \langle \rho_{EM} \rangle = 2 \times 10^{71} \text{ GeV}^4$$

observational value (CMB, ...)

$$\rho_{obs} = 10^{-47} \text{ GeV}^4$$

No solution exists to solve this astounding discrepancy.